

Asymptotic convergence study of a Partial Integro-Differential Equation (PIDE) used to model gene regulatory networks.

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- 1 Introduction
- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
 - Exponential convergence (1D)
 - Exponential convergence evidence (nD)
- 5 Conclusions
- 6 References

Index

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- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
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- 5 Conclusions
- 6 References

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- We derive the partial integral differential (PIDE) model, proposed by Friedman et al., as the continuous counterpart of one master equation with jump processes.



Friedman, N., Cai, L., and Xie, X. S. (2006).

Linking stochastic dynamics to population distribution: An analytical framework of gene expression.

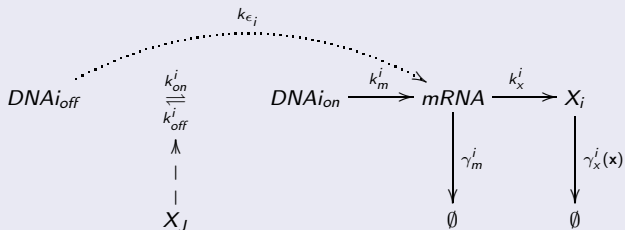
Phys. Rev. Lett., 97(16), 168302.

- Study of self regulated gene expression networks usually involve low copy numbers. Stochastic processes. Chemical Master Equation (CME), its solution is not available in the most cases.
- We derive the partial integral differential (PIDE) model, proposed by Friedman et al., as the continuous counterpart of one master equation with jump processes.
- Using entropy methods we study the convergence to equilibrium, we prove (1D PIDE) or find numerical evidences (nD PIDE) of exponential stability.

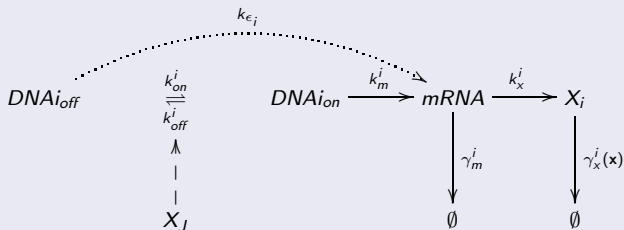
Index

- 1 Introduction
- 2 System description
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Gene Regulatory Network



Gene Regulatory Network



Reaction Steps

1. $\emptyset \xrightarrow{k_m^i c_i(\mathbf{x})} mRNA_i$
2. $mRNA_i \xrightarrow{k_x^i} mRNA_i + X_i$
3. $mRNA_i \xrightarrow{\gamma_m^i} \emptyset$
4. $X_i \xrightarrow{\gamma_x^i(\mathbf{x})} \emptyset$

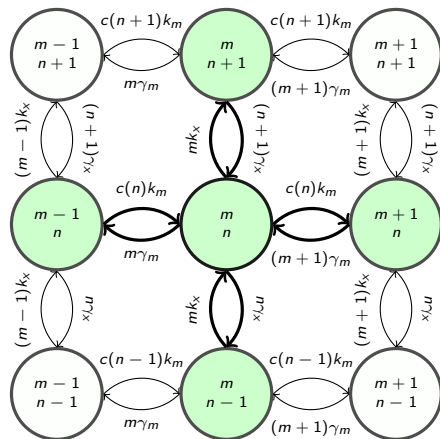
- Self regulation mechanism: $c(x) = [1 - \rho(x)] + \rho(x)\varepsilon$, with $\varepsilon = \frac{k_{\varepsilon}}{k_m} \in (0, 1)$ the transcriptional leakage constant and $\rho(x) = \frac{x^H}{x^H + K^H}$ the Hill type function, where $K = \frac{k_{off}}{k_{on}}$ is an equilibrium constant and H the Hill coefficient.

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- Protein production in bursts: $\omega(x - y) = \frac{1}{b} \exp \left[\frac{-(x - y)}{b} \right]$ the conditional probability for protein level to jump from a state y to x .

Index

- 1 Introduction
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Chemical Master Equation (CME)



$$\begin{aligned} \frac{\partial p_{mn}}{\partial t} = & k_m c(n) (\mathbb{E}_m^{-1} - 1) p_{mn} \\ & + \gamma_m (\mathbb{E}_m^1 - 1) m p_{mn} \\ & + k_x m (\mathbb{E}_n^{-1} - 1) p_{mn} \\ & + \gamma_x (\mathbb{E}_n^1 - 1) n p_{mn} \end{aligned}$$

with \mathbb{E}_m and \mathbb{E}_n being step operators such that:

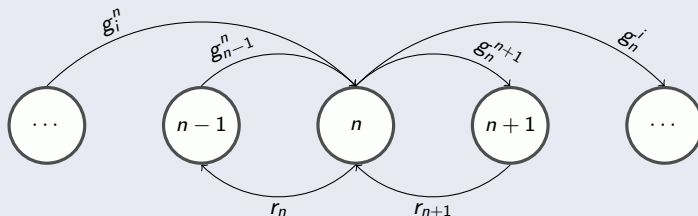
$$\mathbb{E}_m^{-1} p_{mn} = p_{m-1n}$$

$$\mathbb{E}_m^1 m p_{mn} = (m+1) p_{m+1n}$$

$$\mathbb{E}_n^{-1} p_{mn} = p_{mn-1}$$

$$\mathbb{E}_n^1 n p_{mn} = (n+1) p_{mn+1}$$

Master equation deduction

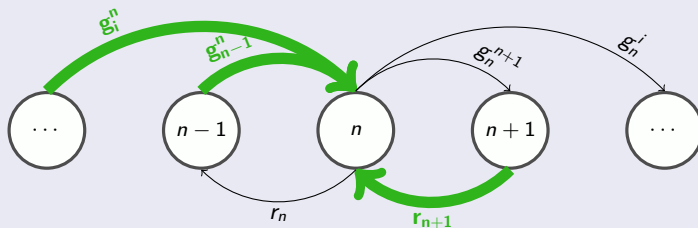


Transition probabilities

- $g_i^n := k_m c(i) \omega(n-i)$
- $r_n := \gamma_x n$

$$P(t + \Delta t, n) = \quad (1)$$

Master equation deduction

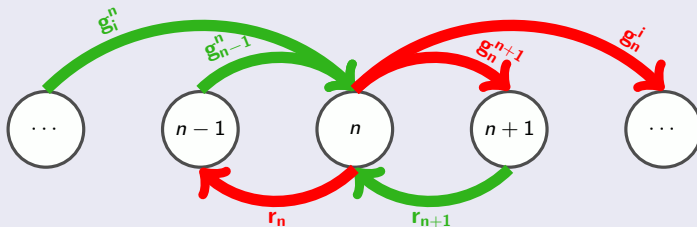


Transition probabilities

- $g_i^n := k_m c(i) \omega(n-i)$
- $r_n := \gamma_x n$

$$P(t + \Delta t, n) = \sum_{i=0}^{n-1} g_i^n P(t, i) \Delta t + r_{n+1} P(t, n+1) \Delta t \quad (1)$$

Master equation deduction

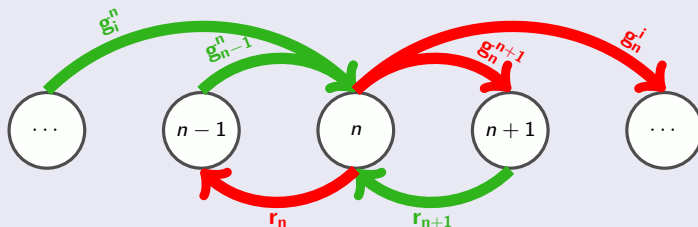


Transition probabilities

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$$P(t + \Delta t, n) = \sum_{i=0}^{n-1} g_i^n P(t, i) \Delta t + r_{n+1} P(t, n+1) \Delta t + P(t, n) \left(1 - r_n \Delta t - \sum_{i=n+1}^{\infty} g_n^i \Delta t \right) \quad (1)$$

Master equation deduction



Transition probabilities

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Master equation with jump processes

$$\frac{\partial P(t, n)}{\partial t} = \sum_{i=0}^{n-1} g_i^n P(t, i) - \sum_{i=n+1}^{\infty} g_n^i P(t, n) + r_{n+1} P(t, n+1) - r_n P(t, n) \quad (2)$$

Continuous formulation (PIDE)

Master equation with jump processes

$$\frac{\partial P(t, n)}{\partial t} = \sum_{i=0}^n g_i^n P(t, i) - \sum_{i=n}^{\infty} g_n^i P(t, n) + r_{n+1} P(t, n+1) - r_n P(t, n) \quad (3)$$

Continuous formulation (PIDE)

Master equation with jump processes

$$\frac{\partial P(t, n)}{\partial t} = \underbrace{\sum_{i=0}^n g_i^n P(t, i)}_{\approx k_m \int_0^x \omega(x-y)c(y)P(t, y)dy} - \sum_{i=n}^{\infty} g_n^i P(t, n) + r_{n+1}P(t, n+1) - r_n P(t, n) \quad (3)$$

The integer indexes n and i are substituted by real x and y respectively:

$$\sum_{i=0}^n g_i^n P(t, i) \approx \int_0^x g_y^x p(t, y) dy = k_m \int_0^x \omega(x-y)c(y)P(t, y)dy$$

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- $\sum_{i=n}^{\infty} g_n^i P(t, n) \approx \int_x^{\infty} g_x^y p(t, x) dy = k_m c(x) p(t, x) \int_x^{\infty} \omega(y-x) dy = k_m c(x) p(t, x)$

Continuous formulation (PIDE)

Master equation with jump processes

$$\frac{\partial P(t, n)}{\partial t} = \sum_{i=0}^n g_i^n P(t, i) - \sum_{i=n}^{\infty} g_n^i P(t, n) + \underbrace{r_{n+1}P(t, n+1) - r_n P(t, n)}_{\approx \frac{\partial}{\partial x} [\gamma_x x P(t, x)]} \quad (3)$$

The integer indexes n and i are substituted by real x and y respectively:

- $\sum_{i=0}^n g_i^n P(t, i) \approx \int_0^x g_y^x p(t, y) dy = k_m \int_0^x \omega(x-y) c(y) P(t, y) dy$
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- $r_{x+1} p(t, x+1) \approx r_x p(t, x) + \frac{\partial [r_x p(t, x)]}{\partial x}$ (Taylor theorem)

Continuous formulation (PIDE)

Master equation with jump processes

$$\frac{\partial P(t, n)}{\partial t} = \sum_{i=0}^n g_i^n P(t, i) - \sum_{i=n}^{\infty} g_n^i P(t, n) + r_{n+1} P(t, n+1) - r_n P(t, n) \quad (3)$$

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- $r_{x+1}p(t, x+1) \approx r_x p(t, x) + \frac{\partial [r_x p(t, x)]}{\partial x}$ (Taylor theorem)

Partial Integral Differential Equation (PIDE)

$$\frac{\partial p(\tau, x)}{\partial \tau} = \frac{\partial [xp(\tau, x)]}{\partial x} - ac(x)p(\tau, x) + a \int_0^x \omega(x-y)c(y)p(\tau, y) dy \quad (4)$$

with dimensionless time, $\tau = \gamma_x t$ and $a = k_m / \gamma_x$.

Generalized PIDE

Generalized Partial Integral Differential Equation

$$\frac{\partial p(t, \mathbf{x})}{\partial t} = \sum_{i=1}^n \left(\frac{\partial}{\partial x_i} \left[\gamma_x^i(\mathbf{x}) x_i p(t, \mathbf{x}) \right] - k_m^i c_i(\mathbf{x}) p(t, \mathbf{x}) + k_m^i \int_0^{x_i} \omega_i(x_i - y_i) c_i(\mathbf{y}_i) p(t, \mathbf{y}_i) dy_i \right)$$

\mathbf{y}_i represents the vector state \mathbf{x} , whose i position is changed for y_i , $((\mathbf{y}_i)_j = x_j \text{ if } j \neq i \text{ and } (\mathbf{y}_i)_j = y_i \text{ if } j = i)$, $\gamma_x^i(\mathbf{x})$ is the degradation rate function of each protein and

$$\omega_i(x_i - y_i) = \frac{1}{b_i} \exp\left(-\frac{x_i - y_i}{b_i}\right)$$

is the conditional probability for protein jumping from a state y_i to x_i after a burst. The function $c_i(\mathbf{x})$ ($c_i : \mathbb{R}_+^n \rightarrow [\varepsilon_i, 1]$) is the input function which models the regulation mechanism.



Pájaro, M., Alonso, A. A., Otero-Muras, I., and Vázquez, C. (2017).

Stochastic Modeling and Numerical Simulation of Gene Regulatory Networks with Protein Bursting.

J. Theor. Biol., 421, 51-70.

Steady state solution PIDE (1D)

Analytic steady state solution for the 1D PIDE

$$P_{\infty}(x) := Z [\rho(x)]^{\frac{a(1-\varepsilon)}{H}} x^{-(1-a\varepsilon)} e^{\frac{-x}{b}}$$

$$= Z \left[x^H + K^H \right]^{\frac{a(\varepsilon-1)}{H}} x^{a-1} e^{\frac{-x}{b}},$$

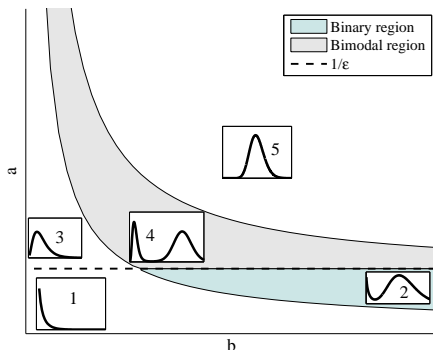
with $\rho(x) = \frac{x^H}{x^H + K^H}$ and Z being a normalizing constant such that $\int_0^{\infty} P_{\infty}(x) dx = 1$.



Pájaro, M., Alonso, A. A., and Vázquez, C. (2015).

Shaping protein distributions in stochastic self-regulated gene expression networks.

Phys. Rev. E, 92(3), 032712.



Index

- 1 Introduction
- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
 - Exponential convergence (1D)
 - Exponential convergence evidence (nD)
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- 6 References

Entropy inequality (I)

General entropy functional

$$\mathcal{G}_H(u) := \int_0^\infty H(u(x)) P_\infty(x) dx, \quad (5)$$

with $H(u)$ being any convex function of $u(x)$ and $u(x) := \frac{p}{p_\infty}(x)$.

Entropy inequality (I)

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Proposition 1

Let $\mathcal{G}_H(u)$ be the general entropy functional, then the following equality is verified

$$\frac{d\mathcal{G}_H(u)}{d\tau} = \mathcal{D}_H(u), \quad (6)$$

$$\mathcal{D}_H(u) = a \int_0^\infty \int_y^\infty \omega(x-y) [H(u(x)) - H(u(y)) + H'(u(x)) (u(y) - u(x))] c(y) P_\infty(y) dx dy.$$



Michel, P., Mischler, S., and Perthame, B. (2005).

General relative entropy inequality: An illustration on growth models.

J. Math. Pures Appl., 84(9), 1235–1260.

Entropy inequality (II)

Entropy functional

We consider $H(u) = (u - 1)^2$ with $u(x) := \frac{p}{P_\infty}(x)$,

$$\mathcal{G}_2(u) := \int_0^\infty \left(\frac{p}{P_\infty}(x) - 1 \right)^2 P_\infty(x) dx = \int_0^\infty u(x)^2 P_\infty(x) dx - 1 \quad (7)$$

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Proposition 1

Let $\mathcal{G}_2(u)$ be the entropy functional, then the following equality is verified

$$\frac{d\mathcal{G}_2(u)}{d\tau} = -\mathcal{D}_2(u), \quad (8)$$

$$\text{with } \mathcal{D}_2(u) = a \int_0^\infty \int_y^\infty \omega(x - y) (u(x) - u(y))^2 c(y) P_\infty(y) dx dy.$$

Numerical evidence of exponential stability

Purpose

To prove that: $\mathcal{G}_2(u) \leq \frac{1}{2\beta} \mathcal{D}_2(u)$

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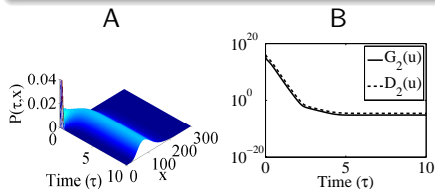


Figure: The temporal evolution of (4) is represented in plot A for parameters $a = 10$, $b = 20$, $H = 1$, $K = 70$ and $\varepsilon = 0.05$. The dashed and continuous lines in plot B represent $\mathcal{G}_2(u)$ and $\mathcal{D}_2(u)$ respectively.

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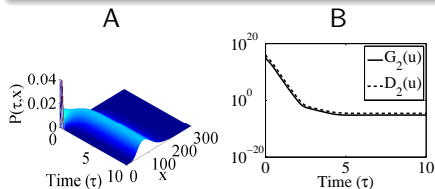


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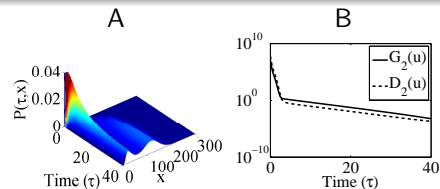


Figure: The temporal evolution of (4) is represented in plot A for parameters $a = 27$, $b = 5$, $H = -4$, $K = 70$ and $\varepsilon = 0.2$. The dashed and continuous lines in plot B represent $\mathcal{G}_2(u)$ and $\mathcal{D}_2(u)$ respectively.

Index

- 1 Introduction
- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
 - Exponential convergence (1D)
 - Exponential convergence evidence (nD)
- 5 Conclusions
- 6 References

Preliminary results (I)

Lemma 1

Let $P_\infty : (0, \infty) \rightarrow \mathbb{R}_+$ be the steady state solution of the 1D PIDE model such that $\int_0^\infty P_\infty(x) dx = 1$. Defining

$$\mathcal{H}_2(u) := \int_0^\infty \int_y^\infty P_\infty(x) P_\infty(y) (u(x) - u(y))^2 dx dy, \quad (9)$$

with $u(x) = \frac{p(x)}{P_\infty(x)}$, there holds:

$$\mathcal{G}_2(u) = \mathcal{H}_2(u). \quad (10)$$

Preliminary results (I)

Lemma 1

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with $u(x) = \frac{p(x)}{P_\infty(x)}$, there holds:

$$\mathcal{G}_2(u) = \mathcal{H}_2(u). \quad (10)$$

$$\mathcal{G}_2(u) := \int_0^\infty \left(\frac{p}{P_\infty}(x) - 1 \right)^2 P_\infty(x) dx = \int_0^\infty u(x)^2 P_\infty(x) dx - 1$$

Preliminary results (II)

Lemma 2 (P_∞ bounds)

For $\delta > 0$ we define the intervals of length $\frac{1}{2}$:

$$I_{k,\delta} := \left(\delta + \frac{k}{2}, \delta + \frac{k+1}{2} \right], \quad k \geq 0 \text{ integer}, \quad (11)$$

and

$$p_k := Z \left[\left(\delta + \frac{k}{2} \right)^H + K^H \right]^{\frac{a(\varepsilon-1)}{H}} \left(\delta + \frac{k}{2} \right)^{a-1} e^{\frac{-(\delta+\frac{k}{2})}{b}} = P_\infty \left(\delta + \frac{k}{2} \right). \quad (12)$$

Then, the following inequality holds:

$$A(\delta) \leq \frac{P_\infty(x)}{p_k} \leq B(\delta), \quad \forall x \in I_{k,\delta} \text{ and } \forall k, \quad (13)$$

with $P_\infty(x) = Z [x^H + K^H]^{\frac{a(\varepsilon-1)}{H}} x^{a-1} e^{\frac{-x}{b}}$ and Z being a normalizing constant such that $\int_0^\infty P_\infty(x) = 1$.

Preliminary results (II)

Proof of Lemma 2 (P_∞ bounds)

$$A(\delta, k) := \begin{cases} \left(\frac{(\delta + \frac{k+1}{2})^H + K^H}{(\delta + \frac{k}{2})^H + K^H} \right)^{\frac{a(\varepsilon-1)}{H}} e^{\frac{-1}{2b}} & \text{if } a \geq 1 \\ \left(\frac{(\delta + \frac{k+1}{2})^H + K^H}{(\delta + \frac{k}{2})^H + K^H} \right)^{\frac{a(\varepsilon-1)}{H}} e^{\frac{-1}{2b}} \left(\frac{2\delta + k + 1}{2\delta + k} \right)^{a-1} & \text{if } a < 1 \end{cases}$$

and

$$B(\delta, k) := \begin{cases} \left(\frac{2\delta + k + 1}{2\delta + k} \right)^{a-1} & \text{if } a > 1 \\ 1 & \text{if } a \leq 1 \end{cases}$$

Notice that,

$$\lim_{k \rightarrow \infty} A(\delta, k) = e^{-\frac{1}{2b}}, \quad \lim_{k \rightarrow \infty} B(\delta, k) = 1$$

and they are $A(\delta) := \min_{k \geq 0} (A(\delta, k))$ and $B(\delta) := \max_{k \geq 0} (B(\delta, k))$

Preliminary results (III)

Lemma 3

We define the term M_j as:

$$M_j := \sum_{k=1}^{j-1} \frac{1}{m_k} \quad (14)$$

with $\{m_k\}_{k \geq 1}$ a positive sequence given by $m_k = p_k e^{\frac{\delta + \frac{k}{2}}{2b}}$. Then, the following inequality is verified:

$$m_k \sum_{j=k+1}^{\infty} M_j p_j \leq C p_k, \quad (15)$$

for some constant $C > 0$.

Preliminary results (III)

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We define the term M_j as:

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$$m_k \sum_{j=k+1}^{\infty} M_j p_j \leq C p_k, \quad (15)$$

for some constant $C > 0$.

$$p_k = P_{\infty} \left(\delta + \frac{k}{2} \right) = Z \left[\left(\delta + \frac{k}{2} \right)^H + K^H \right]^{\frac{a(\varepsilon-1)}{H}} \left(\delta + \frac{k}{2} \right)^{a-1} e^{\frac{-(\delta + \frac{k}{2})}{b}} = P_{\infty} \left(\delta + \frac{k}{2} \right).$$

Preliminary results (III)

Lemma 3

Defining $\{a_j\}_{j \geq 1}$ with $a_j = \frac{1}{m_j}$ we have that

$$\lim_{j \rightarrow \infty} \frac{a_{j+1} - a_j}{M_{j+1} - M_j} = e^{\frac{1}{4b}} - 1$$

Since this limit exist and $\{M_j\}_{j \geq 1}$ is a strictly increasing and divergent sequence, we can use the Stolz-Cesàro theorem to obtain that $M_j \leq C_0 a_j$, with $C_0 > 0$ constant. Then,

$$m_k \sum_{j=k+1}^{\infty} M_j p_j \leq C_0 m_k \sum_{j=k+1}^{\infty} a_j p_j, \quad \text{with} \quad \sum_{j=k+1}^{\infty} a_j p_j = \sum_{j=k+1}^{\infty} e^{-\frac{2\delta+j}{4b}} = \frac{e^{-\frac{2b-1}{4b}}}{e-1} e^{-\frac{2\delta+k}{4b}}$$

So that,

$$m_k \sum_{j=k+1}^{\infty} M_j p_j \leq C m_k e^{-\frac{2\delta+k}{4b}} = C p_k, \quad \text{with} \quad C = C_0 \frac{e^{-\frac{2b-1}{4b}}}{e-1}$$

Main results (I)

Purpose

To prove that: $\mathcal{G}_2(u) = \mathcal{H}_2(u) \leq \frac{1}{2\beta} \mathcal{D}_2(u)$

$$\text{with } \mathcal{D}_2(u) = a \int_0^\infty \int_y^\infty \omega(x-y) (u(x) - u(y))^2 c(y) P_\infty(y) dx dy$$

$$\text{and } \mathcal{H}_2(u) = \int_0^\infty \int_y^\infty P_\infty(x) P_\infty(y) (u(x) - u(y))^2 dx dy$$

Main results (I)

Purpose

To prove that: $\mathcal{G}_2(u) = \mathcal{H}_2(u) \leq \frac{1}{2\beta} \mathcal{D}_2(u)$

$$\text{with } \mathcal{D}_2(u) = a \int_0^\infty \int_y^\infty \omega(x-y) (u(x) - u(y))^2 c(y) P_\infty(y) dx dy$$

$$\text{and } \mathcal{H}_2(u) = \int_0^\infty \int_y^\infty P_\infty(x) P_\infty(y) (u(x) - u(y))^2 dx dy$$

Proposition 2

The following inequality is verified:

$$\lambda \mathcal{H}_2(u) \leq \int_0^\infty \int_y^{y+1} P_\infty(y) (u(x) - u(y))^2 dx dy := D(u), \quad (16)$$

for some constant $\lambda > 0$.

Main results (II)

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The following inequality is verified:

$$\lambda \mathcal{H}_2(u) \leq \int_0^\infty \int_y^{y+1} P_\infty(y) (u(x) - u(y))^2 dx dy, \quad \text{for some constant } \lambda > 0.$$

Proposition 3

The following inequality is verified:

$$\alpha \int_0^\infty \int_y^{y+1} P_\infty(y) (u(x) - u(y))^2 dx dy \leq \mathcal{D}_2(u) \quad (17)$$

for some constant $\alpha > 0$. As consequence, we deduce the exponential convergence towards P_∞ .

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Main results (Proof) (I)

Proof of Proposition 2

We take $0 < \delta < 1$ and split $\mathcal{H}_2(u)$ in two parts:

$$\begin{aligned}\mathcal{H}_2(u) &= \int_{\delta}^{\infty} \int_y^{\infty} P_{\infty}(x) P_{\infty}(y) (u(x) - u(y))^2 dx dy \\ &\quad + \int_0^{\delta} \int_y^{\infty} P_{\infty}(x) P_{\infty}(y) (u(x) - u(y))^2 dx dy := \mathcal{H}_{21}(u) + \mathcal{H}_{22}(u)\end{aligned}$$

For $i, j \geq 0$ integers we define:

$$A_{i,j} := \int_{I_{i,\delta}} \int_{I_{j,\delta}} (u(x) - u(y))^2 dy dx = \int_{I_{i,\delta}} \int_{I_{j,\delta}} (u(x) - u(y))^2 dx dy.$$

We can estimate both the left and the right hand side of (16) by using the quantities $A_{i,j}$.

Main results (Proof) (II)

Proof of Proposition 2

$$\mathcal{H}_{21}(u) \leq B(\delta)^2 \sum_{i=0}^{\infty} \sum_{j=0}^i p_i p_j A_{i,j} \quad (18)$$

$$\sum_{i=0}^{\infty} p_i^2 A_{i,i} \leq \frac{P_M}{A(\delta)^2} D(u) \quad (19)$$

Using that $A_{i,j} \leq M_j \sum_{k=i}^{j-1} m_k A_{k,k+1}$ for all $j > i$, we have

$$\begin{aligned} \sum_{i=0}^{\infty} \sum_{j=i+1}^{\infty} p_i p_j A_{i,j} &\leq \sum_{i=0}^{\infty} \sum_{j=i+1}^{\infty} p_i p_j M_j \sum_{k=i}^{j-1} m_k A_{k,k+1} = \sum_{k=0}^{\infty} m_k A_{k,k+1} \sum_{j=k+1}^{\infty} p_j M_j \sum_{i=0}^k p_i \\ &\leq C_{\delta}^1 \sum_{k=0}^{\infty} A_{k,k+1} m_k \sum_{j=k+1}^{\infty} M_j p_j \leq C \sum_{k=0}^{\infty} A_{k,k+1} p_k \leq \frac{C}{A(\delta)} D(u) \end{aligned} \quad (20)$$

$$\lambda_1 \mathcal{H}_{21}(u) \leq D(u) \quad (21)$$

Index

1 Introduction

2 System description

3 From CME to PIDE

4 Convergence to equilibrium

- Exponential convergence (1D)

- Exponential convergence evidence (nD)

5 Conclusions

6 References

Entropy inequality nD (I)

General entropy functional

$$\mathcal{G}_H^n(u) = \int_{\mathbb{R}_+^n} H(u(\mathbf{x})) P_\infty(\mathbf{x}) d\mathbf{x}, \quad (22)$$

with $H(u)$ being any convex function of $u(\mathbf{x})$ and $u(\mathbf{x}) := \frac{p}{p_\infty}(\mathbf{x})$.

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Proposition 4

Let $\mathcal{G}_H^n(u)$ be the general entropy functional, then the following equality is verified

$$\frac{d\mathcal{G}_H(u)}{d\tau} = \mathcal{D}_H^n(u), \quad (23)$$

$$\begin{aligned} \mathcal{D}_H^n(u) = & \sum_{i=1}^n k_m^i \int_{\mathbb{R}_+^n} \int_{y_i}^{\infty} \omega_i(x_i - y_i) [H(u(\mathbf{x})) - H(u(\mathbf{y}_i)) + H'(u(\mathbf{x}))(u(\mathbf{y}_i) - u(\mathbf{x}))] \\ & \times c_i(\mathbf{y}_i) P_\infty(\mathbf{y}_i) d\mathbf{x}_i d\mathbf{y}_i \end{aligned}$$

where \mathbf{y}_i represents the vector state \mathbf{x} , whose i component is replaced by y_i .

Entropy inequality nD (II)

Entropy functional

We consider $H(u) = (u - 1)^2$ with $u(\mathbf{x}) := \frac{p}{P_\infty}(\mathbf{x})$,

$$\mathcal{G}_2^n(u) := \int_{\mathbb{R}_+^n} \left(\frac{p}{P_\infty} - 1 \right)^2 P_\infty d\mathbf{x} = \int_{\mathbb{R}_+^n} u^2 P_\infty d\mathbf{x} - 1 \quad (24)$$

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Proposition 1

Let $\mathcal{G}_2^n(u)$ be the entropy functional, then the following equality is verified

$$\frac{d\mathcal{G}_2^n(u)}{d\tau} = -\mathcal{D}_2^n(u), \quad (25)$$

$$\text{with } \mathcal{D}_2^n(u) = \sum_{i=1}^n -k_m^i \int_{\mathbb{R}_+^n} \int_{y_i}^\infty \omega_i(x_i - y_i) [u(\mathbf{x}) - u(\mathbf{y}_i)]^2 c_i(\mathbf{y}_i) P_\infty(\mathbf{y}_i) d\mathbf{x}_i d\mathbf{y}_i$$

Numerical evidence of exponential stability 2D

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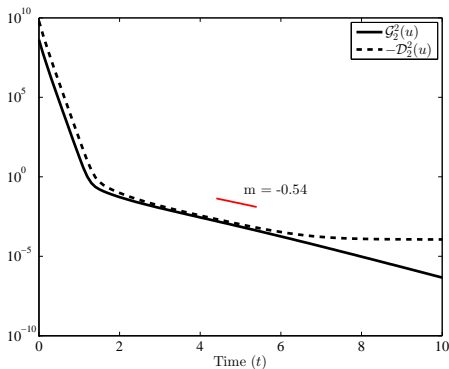
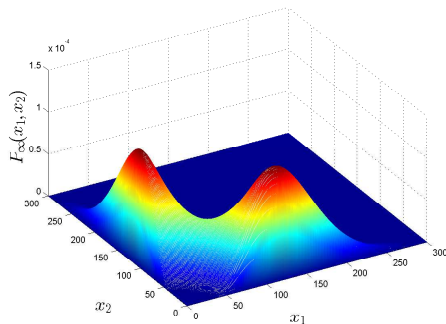


Figure: Steady state distribution of a mutual repression example and time evolution of $\mathcal{G}_2^n(u)$ and $-\mathcal{D}_2^n(u)$.

Index

- 1 Introduction
- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
 - Exponential convergence (1D)
 - Exponential convergence evidence (nD)
- 5 Conclusions**
- 6 References

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Index

- 1 Introduction
- 2 System description
- 3 From CME to PIDE
- 4 Convergence to equilibrium
 - Exponential convergence (1D)
 - Exponential convergence evidence (nD)
- 5 Conclusions
- 6 References**

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